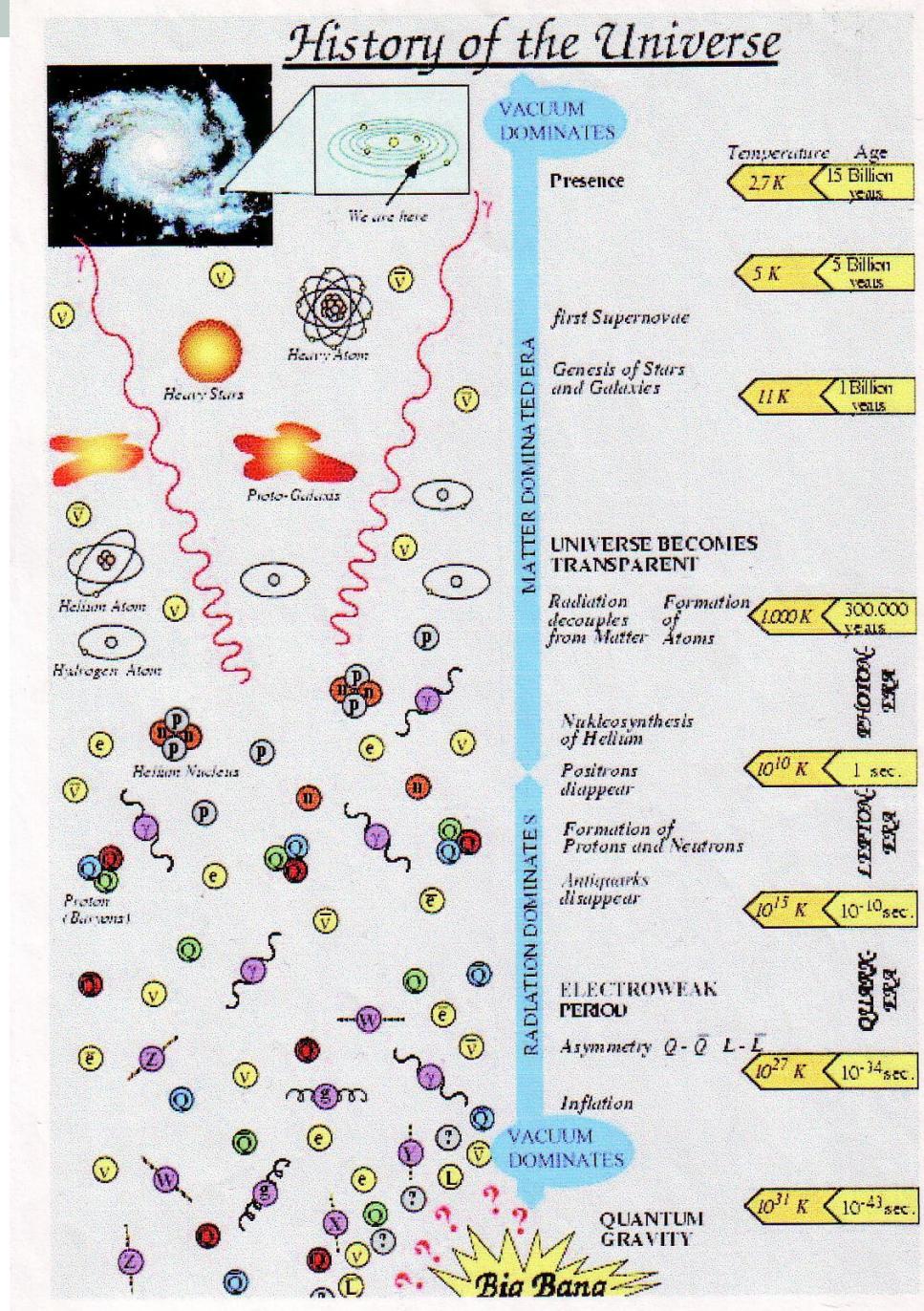
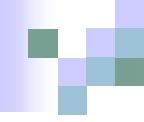


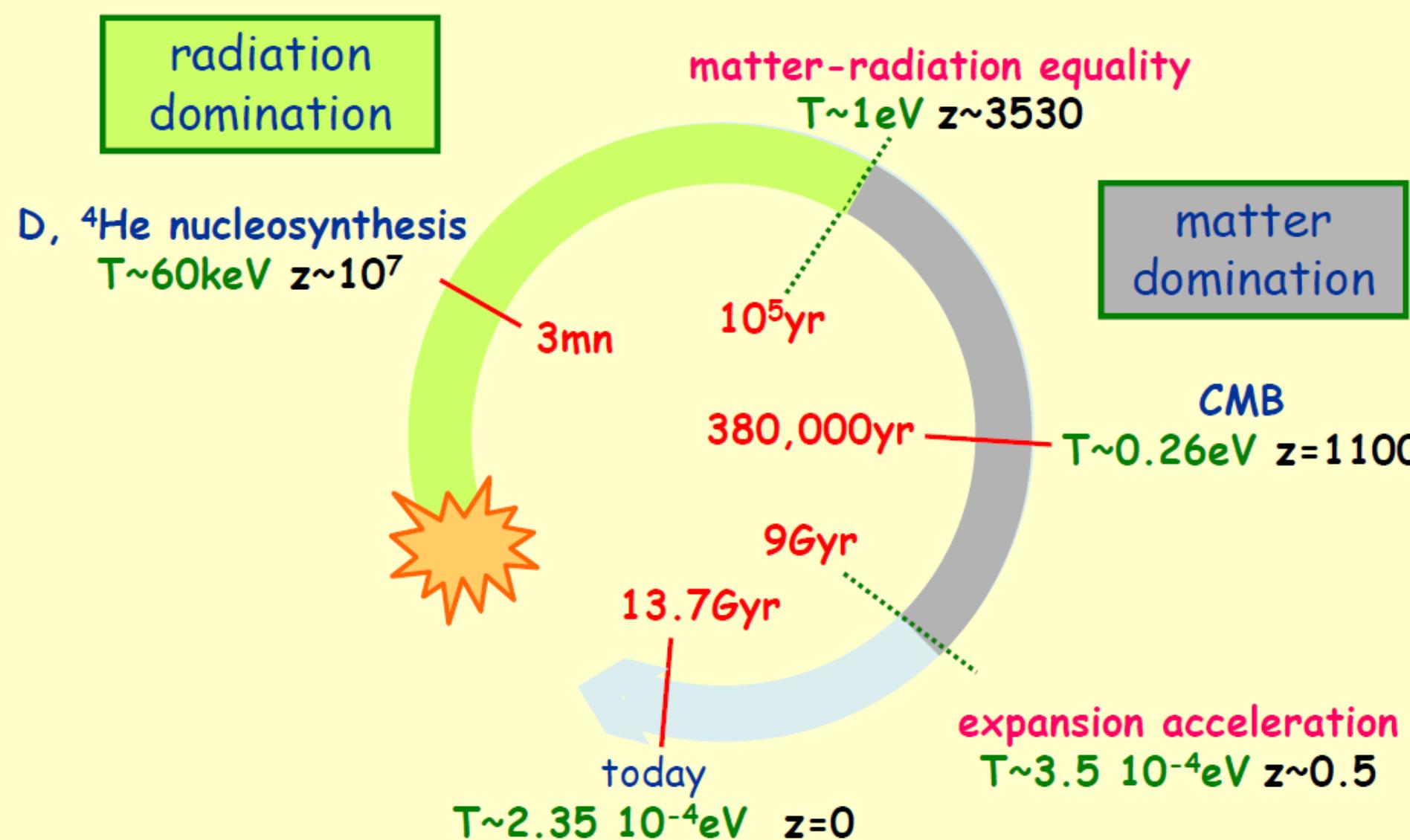
FIZIKALNA KOZMOLOGIJA

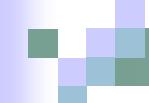
VII. VRLO RANI SLEMIR & INFLACIJA





"KOZMIČKI SAT" ranog svemira





Ekstra zračenje u mjerenuju CMB

$$\rho_r = \rho_\gamma + \rho_\nu \quad \rho_\gamma = \frac{\pi^2}{15} T_\gamma^4,$$

$$\rho_\nu + \rho_{\text{er}} \equiv \frac{\pi^2}{15} T_\nu^4 \cdot \frac{7}{8} N_{\text{eff}}$$

$$\rho_r = \rho_\gamma \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right] \quad N_{\text{eff}} = 3.30 \pm 0.27$$

■ Usporedba s rezultatima LEP-a

$$N_\nu = 2.984 \pm 0.008$$

Usporedba CMB i neutrina

- Vj.: Pozadinsko zračenje neutrina u svjetlu činjenice masivnih neutrina

$$m(\nu_\mu)^2 - m(\nu_e)^2 = 5 \cdot 10^{-5} \text{ eV}^2$$

$$m(\nu_\tau)^2 - m(\nu_\mu)^2 = 3 \cdot 10^{-3} \text{ eV}^2$$

Termička povijest svemira

— u eri zračenja —

Efektivni broj st. slobode $g_* = g_B + \frac{7}{8} g_F$

u ovisnosti o

temperaturi / česticama u ravnoteži

$$k_B T < m_A c^2$$

$$m_S c^2 \quad \gamma \quad e^\pm \quad \nu_e \quad \nu_\tau \quad \mu^\pm \quad \left\{ \begin{array}{l} u\bar{u} \\ d\bar{d}, g \end{array} \right\} \quad \frac{205}{4}$$

$$\downarrow \quad 3 \left\{ \begin{array}{l} \pi^+, \pi^-, \pi^0 \end{array} \right\} \quad \frac{69}{4}$$

$$\boxed{\frac{59}{4}}$$

$$\frac{43}{4}$$

$$\boxed{\frac{11}{2}}$$

jednosmjerno podgrijavajući γ -zrac
 $e^+ e^- \rightarrow \gamma + \gamma \Rightarrow T_\gamma = 1.4 T_B$

2 plota zadajući raspreženje

CMB

3.5 MeV

2.3 MeV

1 MeV

0.2 MeV

1 eV

0.3 eV

VRLO RANI SVEMIR

Vraćanje u prošlost

$\sim 10^5$ god → 10^4 god → 3 min → prije
 (formir. CMB) (dominac zrač.) (dominac ^{4}He) → $10^{-4} \text{ s} / 10^{12} \text{ K}$
 plin mezon, leptoni; γ u tadašnjem.

$\gamma, \nu_e, \nu_\mu, \nu_\tau, e^\pm, \mu^\pm, u\bar{u}, d\bar{d}, g$

$$g_* = \frac{205}{4}$$

$m_S c^2$ - - - - -
 $+ s, \bar{s}$

$$2 \cdot (3 \cdot 2 \cdot \frac{7}{8}) = \frac{42}{4} \Rightarrow \frac{243}{4}$$

$m_e c^2$ - - - - -
 $+ e, \bar{e}$

$$\frac{239}{4}$$

$m_T c^2$ - - - - -
 $+ \tau, \bar{\tau} / b, \bar{b} / W^+, W^-, Z$

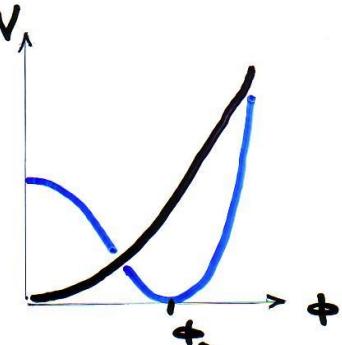
$m_b c^2 / m_{W, Z} c^2$

◆ Fermijova skala SSS

$$V(\phi) = \frac{1}{2} m_\phi^2 \phi^2 + \frac{\lambda}{4} \phi^4$$

$$\langle \phi \rangle = 0 \rightarrow \phi_0 = \pm \sqrt{\frac{\lambda}{2}}$$

$$= 250 \text{ GeV}$$



◆ "GUT" skala

veličina ujednjenja $\sim 10^{15} \text{ GeV} / 10^{-36} \text{ s}$

◆ Planckova skala $\sim 10^{19} \text{ GeV} / 5 \cdot 10^{-44} \text{ s}$

Svemir u vrlo ranoj fazi je BARIONSKI ASIMETRIČAN

- Vj.: Barionsko simetričan svemirna temp. 1 TeV vodio bi na omjer broja nukleona i fotona daleko ispod opaženog

$$\frac{n_N}{n_\gamma} = \frac{g_N \left(\frac{m_N T}{2\pi}\right)^{3/2} e^{-m_N/T}}{2 \frac{1.2022 T^3}{\pi^2}} = 7.3 \times 10^{-27}$$

puno manji od opaženog 6×10^{-10} .

Bariogeneza

- očuvani barionski broj je prije formiranja hadrona (na 200 MeV) nošen kvarkovima

Opužene
gustoće broja $n_B = n_q - n_{\bar{q}}$

→ bezdimenzionalni omjeri :

$$\eta_B = \frac{n_B}{n_\gamma}$$

n_γ = gustoća br. fotona

$$Y_B = \frac{n_B}{S}$$

S = gustoća entropije

Na niskim temperaturama ($T < m_e$)

$$\gamma_B \approx 7 Y_B$$

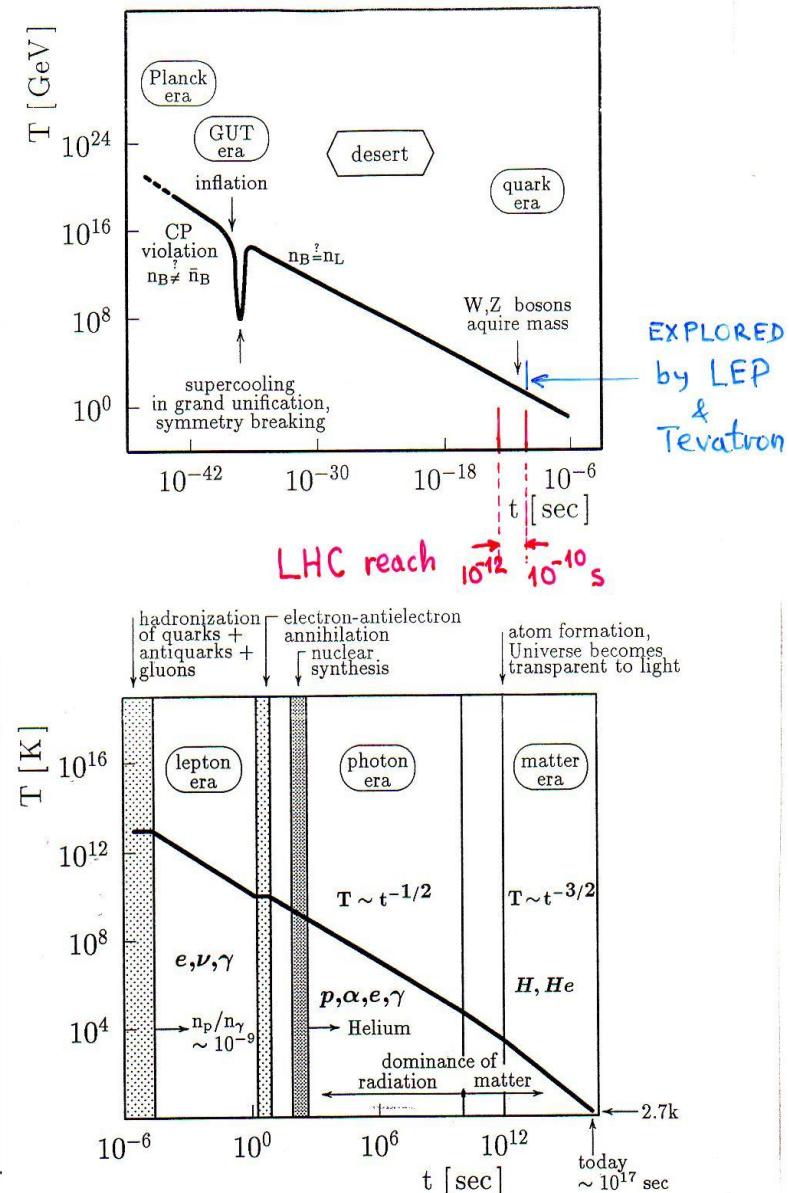
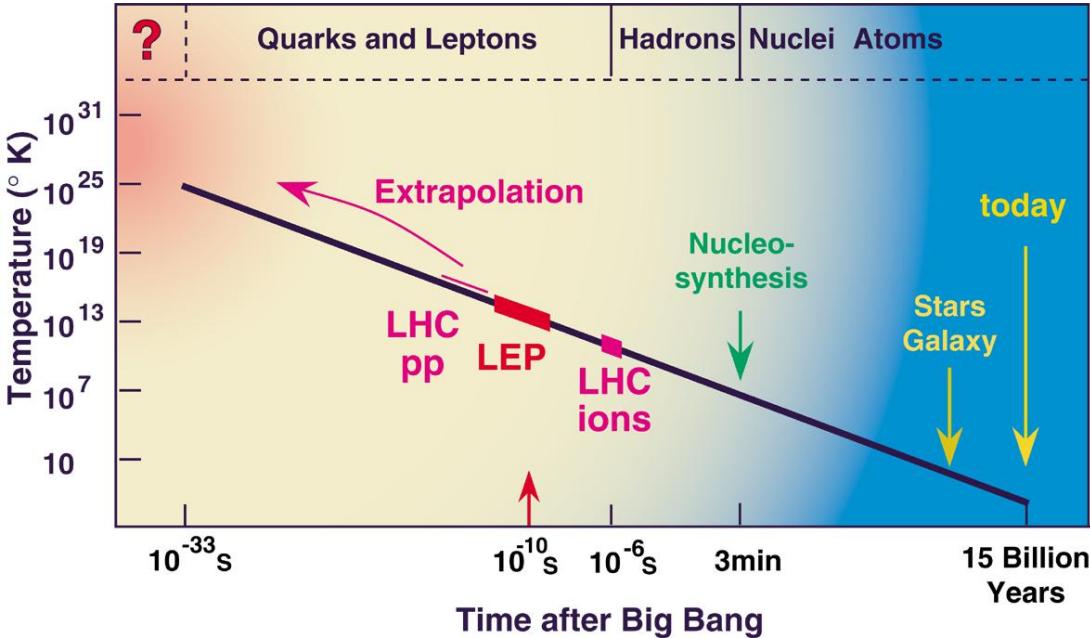
→ mjerene vrijednosti svode se na sadašnje vrijeme (t_0):

$$\gamma_B = 6.1(3) \cdot 10^{-10} \quad \begin{pmatrix} \text{iz CMB-a} \\ \text{WMAP} \end{pmatrix}$$

$$Y_B = 0.8 \cdot 10^{-10}$$

upravo vrijednosti potrebne za nastanak galaktika?

VRLO RANI SVEMIR



NIZ PROBLEMA STANDARDNE (FRW) KOZMOLOGIJE ZAHTIJEVA DA ERI ZRAČENJA PRETHODI INFLATORNA ERA

VRLO RANI SVEMIR

| Vraćanje u prošlost | prije | $g_* = 69/4$ |
|----------------------------------|---|---|
| $\sim 10^5$ god (formir. CMB) | 10^4 god (dominac zvč.) \rightarrow trvar | 3 min (dominantne izgore ^4He) $10^{-4}\text{ s} / 10^{12}\text{ K}$ plin mezonata, (leptona; 1°) |

$$\gamma, \nu_e, \nu_\mu, \nu_\tau, e^\pm, \mu^\pm, u, \bar{u}, d, \bar{d}, g$$

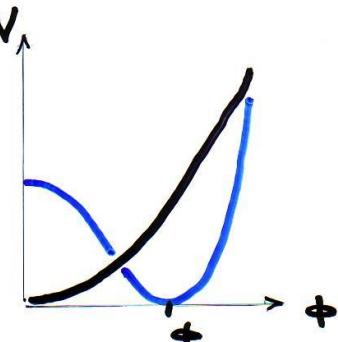
$$\begin{array}{l}
 M_{SC^2} \quad \dots \quad + s, \bar{s} \quad 2 \cdot (3 \cdot 2 \cdot \frac{7}{8}) = \frac{42}{4} \Rightarrow \frac{247}{4} \\
 M_{CC^2} \quad \dots \quad + c, \bar{c} \quad \frac{289}{4} \\
 M_{\tau C^2} \quad \dots \quad + \tau, \bar{\tau} / b, \bar{b} / t, \bar{t}
 \end{array}$$

Fermijova skala

$$V(\phi) = \frac{1}{2} m^2 \phi^2 + \frac{\lambda}{4} \phi^4$$

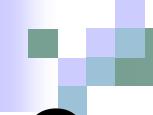
$$\langle \phi \rangle = 0 \rightarrow \phi_0 = \pm \hbar/\sqrt{2}$$

≈ 250 GeV



♦ "GUT" skala
velikost ujednjava s $\sim 10^{15}$ GeV / 10^{-36} s

◆ Planck scale $\sim 10^{19}$ GeV / $5 \cdot 10^{-44}$ s



0-ti PROBLEM: **PROBLEM EKSPANZIJE SVEMIRA** **(ekspanzija sada, jer je bila i u prošlosti) - 2 pristupa:**

- za sve probleme vezane uz “početne uvjete svemira”, kriva je kvantna gravitacija
- kozmologija 80-tih potražila je odgovore na “GUT skali” – na tragu problema magnetnih monopola

Spuštanje sa skale

KVANTNE GRAVITACIJE $\sim 10^{19}$ GeV
na skalu

VELIKE UNIFIKACIJE , GU \bar{I} $\sim 10^{15} - 10^{16}$
GeV

vodi na opću sliku VRLO RANOG
SVEMIRA , koju treba provjeriti opažanjima ,
primjerice PROBLEM MAGNETNIH

MONOPOLA { mase $\gtrsim 10^{-9}$ gr
veličine $\sim 10^{-29}$ cm

kojih bi trebalo biti koliku i fotona CMB-a



◊ "Normalnom" ekspanzijom

jezgre monopola udaljene $\sim 10^{-29}$ cm

na $T_{GUT} \approx 10^{28}$ K

za faktor

$$\frac{T_{GUT}}{T_0 = 3K} = 3 \cdot 10^{27}$$

$$\Rightarrow 3 \cdot 10^{-2}$$

cm danas bi bili dominantni sastojak sremira

◊ Inflacija poreča separaciju

za

$$e^\eta \frac{T_{GUT}}{T_0}$$

$\left\{ \begin{array}{l} \text{za } \eta = 60 \quad 1 \text{ mon/galakt.} \\ \text{za } \eta = 67 \quad 1 \text{ mon/Hubble} \\ \qquad \qquad \qquad \text{sfera} \end{array} \right.$

Koncept inflacije

The idea (A. Guth and A. Linde, 1981): Shortly after the Big Bang, the Universe went through a phase of rapid (exponential) expansion. In this phase the energy and thus the dynamics of the Universe was determined by a term similar to the cosmological constant (vacuum energy).

Why would the Universe do that ?

Why does it help ?

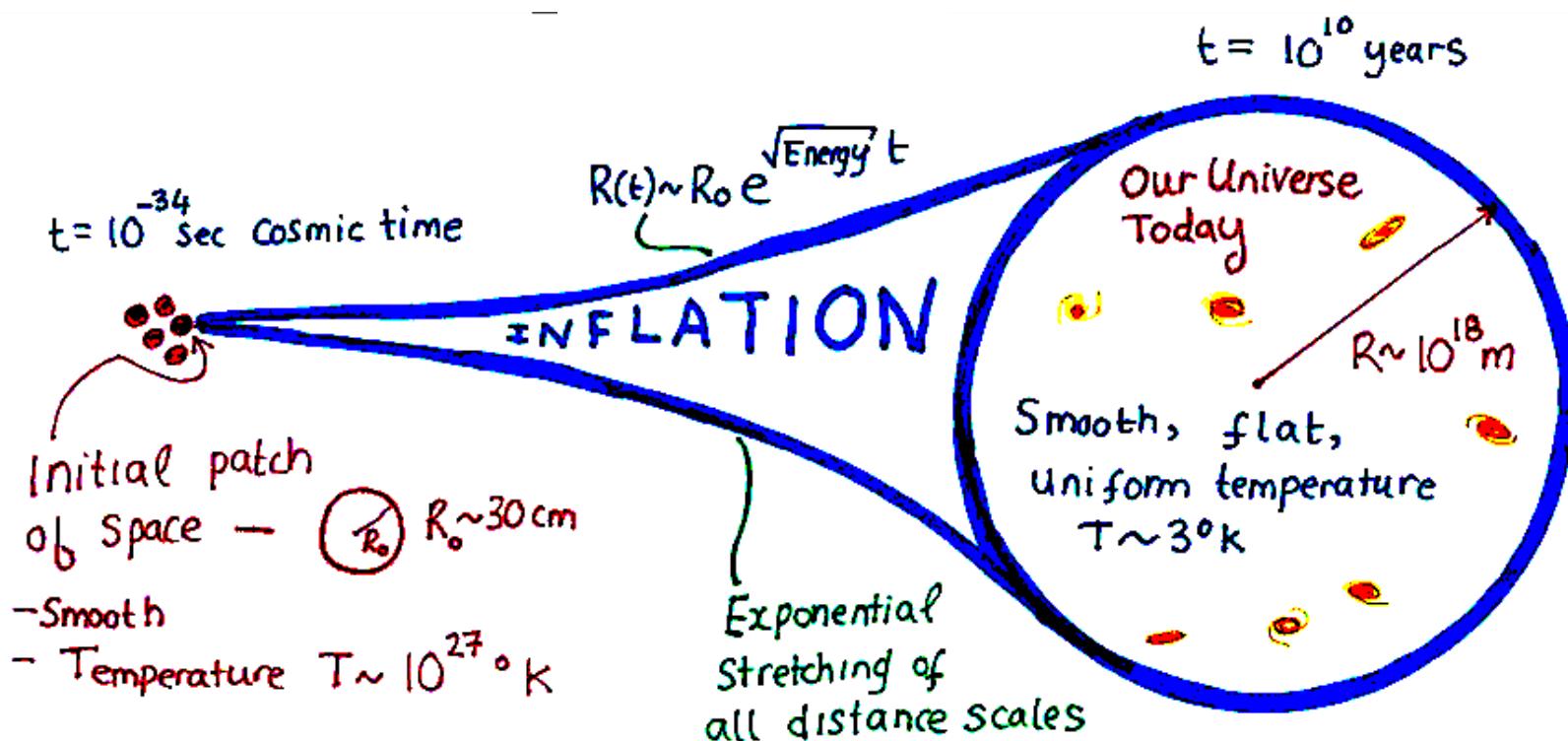
Problem relikta

■ The monopole problem

□ big issue in early 1980s

- Grand Unified Theories of particle physics → at high energies the strong, electromagnetic and weak forces are unified
- the symmetry between strong and electroweak forces 'breaks' at an energy of $\sim 10^{15}$ GeV ($T \sim 10^{28}$ K, $t \sim 10^{-36}$ s)
 - this is a phase transition similar to freezing
 - expect to form 'topological defects' (like defects in crystals)
 - point defects act as magnetic monopoles and have mass $\sim 10^{15}$ GeV/ c^2 (10^{-12} kg)
 - expect one per horizon volume at $t \sim 10^{-36}$ s, i.e. a number density of 10^{82} m $^{-3}$ at 10^{-36} s
 - result: universe today completely dominated by monopoles (not!)

Inflacija i problem relikta

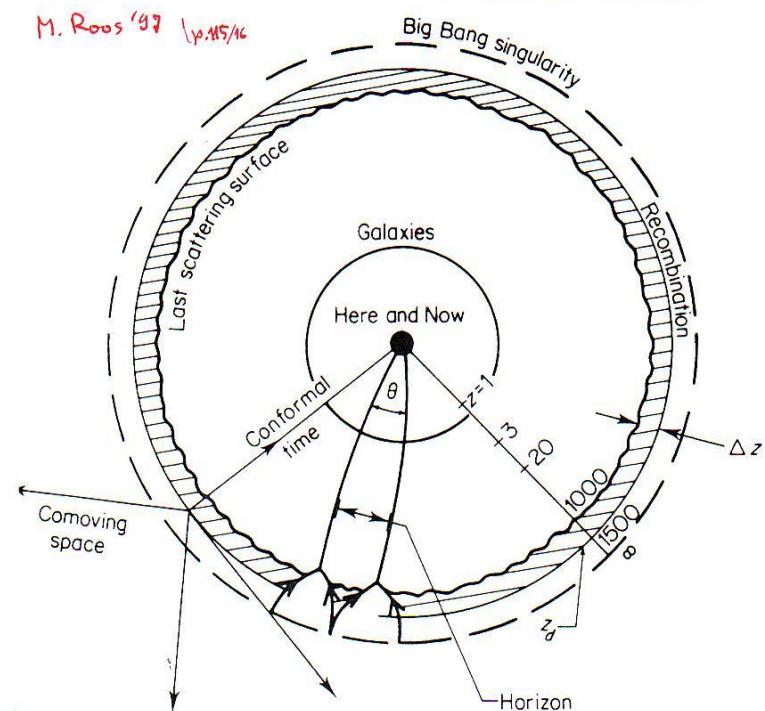


1. PROBLEM: PROBLEM HORIZONTA (područja koja nisu bila u kauzalnom kontaktu nalazimo s istom temperaturom)

Udalj. najudaljenijih objekata u čas t

$$d_{\text{Horiz.}}(t) = S(t) \times (\text{svjetlosni koor. udalj.})$$
$$\int_0^t \frac{c \, dt'}{S(t')}$$
$$\int_0^{t_{\text{odvez}}} \frac{dt'}{S(t')} \ll \int_{t_{\text{odvez}}}^{t_0} \frac{dt'}{S(t')}$$

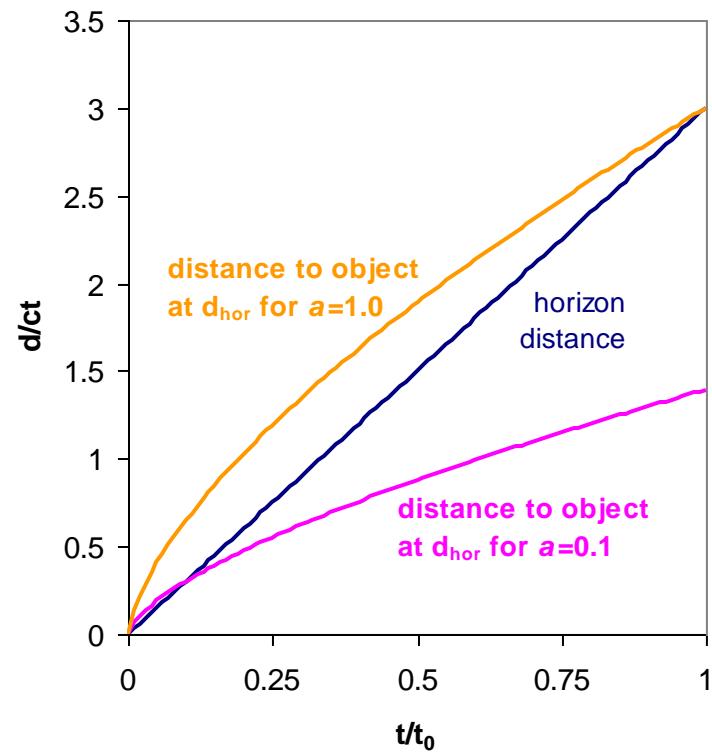
Udaljenost koju je svjetlost možla prevesti prije odvezivanja CMB danas umnožak 1,12° sadarjava horizontu



Horizon Problem

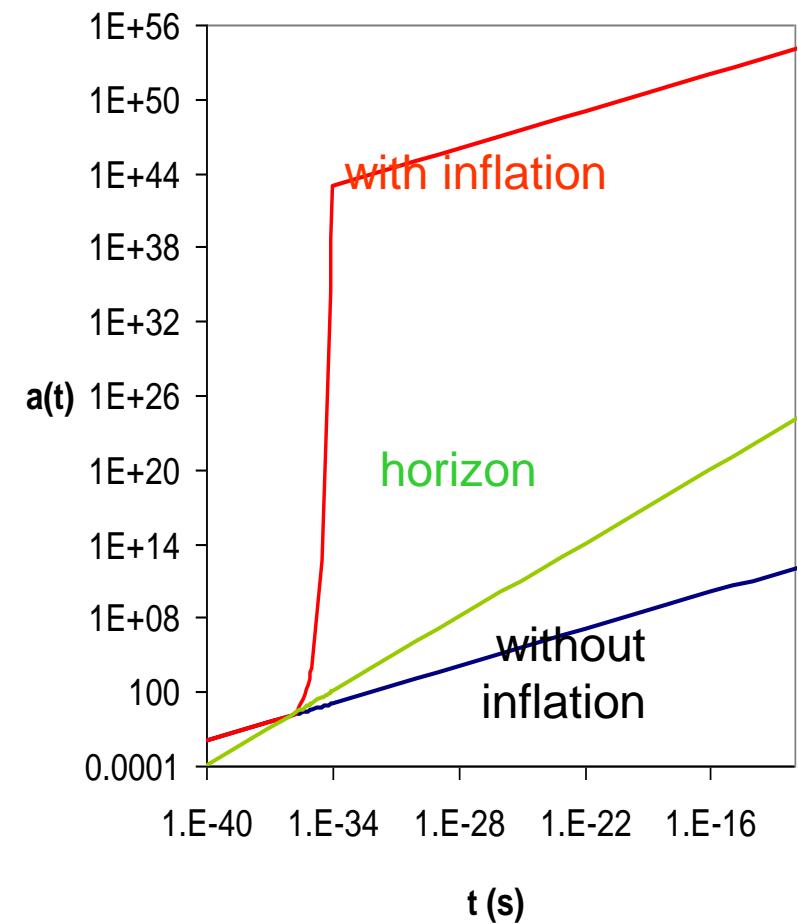
■ Why is the CMB so isotropic?

- consider matter-only universe:
 - horizon distance $d_H(t) = 3ct$
 - scale factor $a(t) = (t/t_0)^{2/3}$
 - therefore horizon expands faster than the universe
 - "new" objects constantly coming into view
- CMB decouples at $1+z \sim 1000$
 - i.e. $t_{CMB} = t_0/10^{4.5}$
 - $d_H(t_{CMB}) = 3ct_0/10^{4.5}$
 - now this has expanded by a factor of 1000 to $3ct_0/10^{1.5}$
 - but horizon distance now is $3ct_0$
 - so angle subtended on sky by one CMB horizon distance is only $10^{-1.5}$ rad $\sim 2^\circ$
- patches of CMB sky $>2^\circ$ apart should not be causally connected



Inflation and the horizon

- Assume large positive cosmological constant Λ acting from t_{inf} to t_{end}
- then for $t_{\text{inf}} < t < t_{\text{end}}$
 $a(t) = a(t_{\text{inf}}) \exp[H_i(t - t_{\text{inf}})]$
 - $H_i = (\frac{1}{3} \Lambda)^{1/2}$
 - if Λ large a can increase by many orders of magnitude in a very short time
- Exponential inflation is the usual assumption but a power law $a = a_{\text{inf}}(t/t_{\text{inf}})^n$ works if $n > 1$



2. PROBLEM: PROBLEM RAVNOSTI

1. predviđanje
inflacije – ravni
svemir
(Friedmannova
jedn. daje de
Sitterovo rješ. za
 $k=0$)

Zanemarimo li "kozmološki član", u okolini $\Omega=1$
Friedmannova j-ba ima oblik

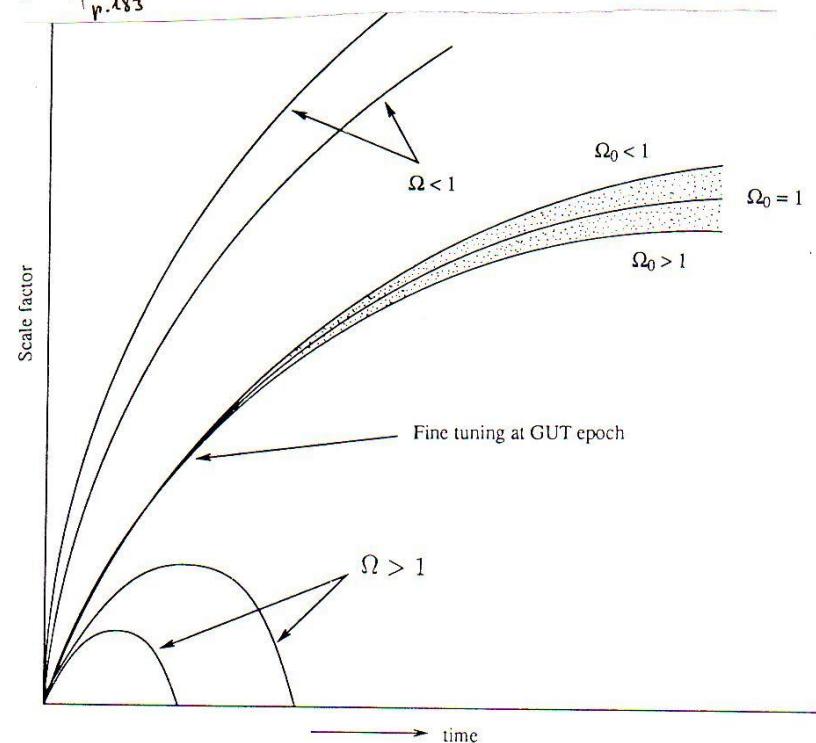
$$|\Omega - 1| = \frac{141}{(S^2 H^2)} \propto \begin{cases} t^{2/3} & \text{dominacija tvori} \\ t^{1/2} & \text{dominacija brčenja} \end{cases}$$

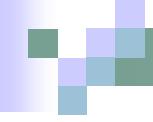
opada
 u stvar. evoluc.
 velikom prostoru

pokazuje da Ω u prošlosti mora biti fini podešen

$$|\Omega - 1| < \begin{cases} \Theta(10^{-16}) & \text{u eri} \\ \Theta(10^{-24}) & \text{nuklearne sinteze} (t \sim 1s) \\ \Theta(10^{-53}) & e-w skale (t \sim 10^{11}s) \\ \Theta(10^{-61}) & GUT skale (t \sim 10^{35}s) \\ & Planck skale (t \sim 10^{-44}s) \end{cases}$$

Nuelic
Fig. 6.7 | p. 183





3. PROBLEM: PROBLEM STRUKTURA

- Problem porijekla svega opaženog!
**Modeli hladne tamne tvari (CDM)
objašnjavaju kako fluktuacije u CMB
rastu do formiranja galaktika**
- Preostaje objasniti odakle fluktuacije
opažene u CMB-u

Inflatorno rješenje problema struktura

- Prije inflacije: postoje kvantne fluktuacije
- Inflacija pojačava kvantne fluktuacije na makroskopske skale
- Nakon inflacije makroskopske fluktuacije (kakve su opažene u CMB zračenju) daju sjeme formiranja galaktika

JOŠ O INFLACIJI

- Da bi očuvali uspjehe modela velikog praska, epohu inflacije treba ograničiti na dovoljno ranu fazu svemira - povratak na normalnu jedn. stanja je kozmološki FAZNI PRIJELAZ
- 2. predviđanje inflacije - pozadinski gravitacijski valovi kao RELIKT dostupan misiji PLANCK-ova mjerena
- Kandidati za pogonitelja inflacije?