Singlet Penguin Contribution to the $B \to K\eta'$ Decay

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[J.O. Eeg, K.K. and I. Picek, Phys. Lett. **B363** (2003) 87]

Overview

- $B \to K\eta'$ decay — experimental data — motivation
- Singlet-penguin as an enhancement mechanism
- $SU(3)_F$ flavour symmetry approach
- Perturbative, dynamical approach
- Conclusions

Experimental data

- CLEO, Belle and BaBar collaborations see a lot of η''s in charmless (rare) hadronic B decays ...

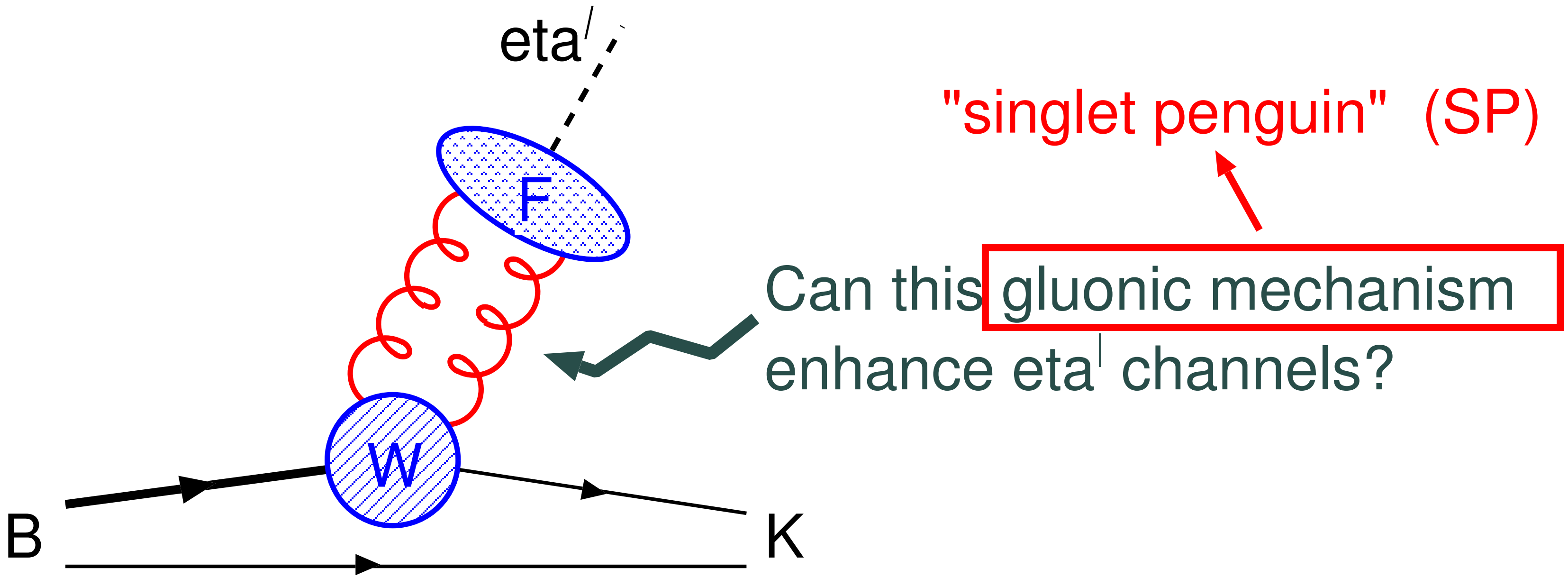
$$\text{Br}(B^+ \to K^+\eta') = (77 \pm 5) \cdot 10^{-6}$$
$$\text{Br}(B^0 \to K^0\eta') = (61 \pm 6) \cdot 10^{-6}$$

- ... as compared to the π's:

$$\text{Br}(B^+ \to K^+\pi^0) = (13 \pm 1) \cdot 10^{-6}$$
$$\text{Br}(B^0 \to K^0\pi^0) = (11 \pm 1) \cdot 10^{-6}$$

- Why are η' channels enhanced?

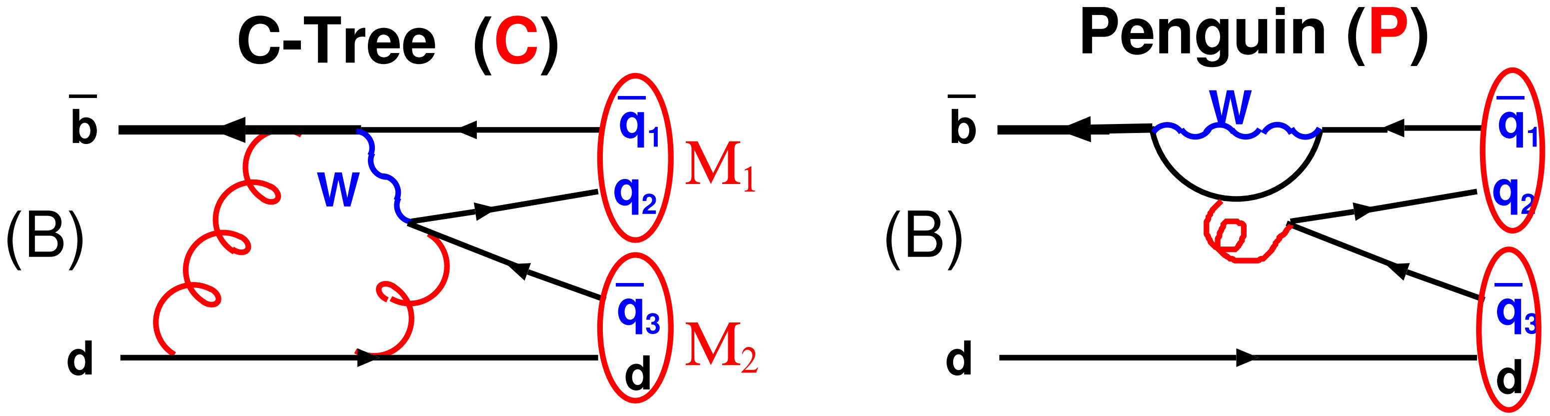
- Experience with η' mass (*U(1) problem*: $m_{\eta'} \gg m_\pi$) suggests: $|\eta'\rangle = \cdots + |gg\rangle + \cdots$



- 1. $SU(3)_F$ symmetry approach $\longrightarrow$ SP part up to 50 %
- 2. perturbative approach $\longrightarrow$ SP part negligible!

$SU(3)_F$ flavour symmetry approach

- decomposing amplitude on various flavour topologies:

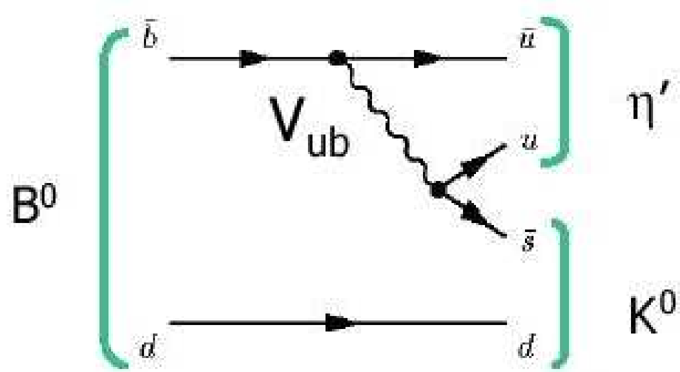


- other topologies: tree (T), exchange (E), annihilation (A), penguin-annihilation (PA), singlet penguin (SP)
- cannot calculate C, T, P, SP, ... but hope that they are invariant under flavour rotations $q_i = u \leftrightarrow d \leftrightarrow s$

SU(3)$_F$ flavour symmetry approach

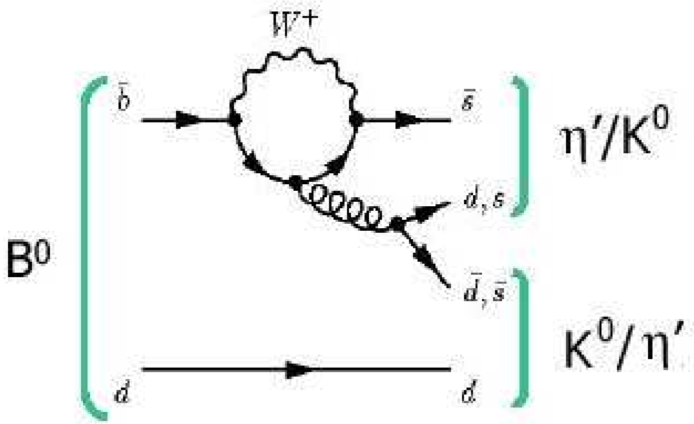
- 7 free parameters — to be predictive one assumes that some can be neglected

C-Tree C $=$



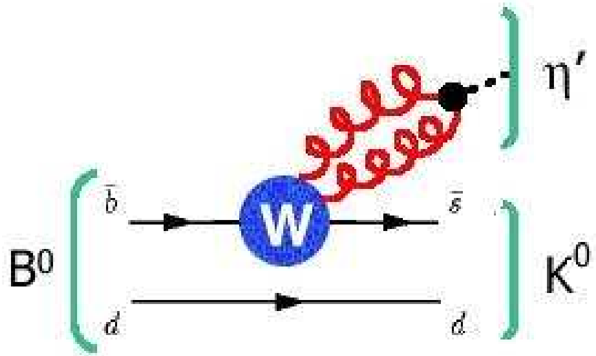
$\longrightarrow$ suppressed by small V_{ub}

Penguin P $=$



$\longrightarrow$ important contribution (interfer-ence!)

Singlet penguin part

- $SP =$ 

- [Chiang, Gronau, Rosner (2003)]: $SP/P \approx 0.4 - 0.8$
- Possible objections:
  - $SU(3)_F$ symmetry broken
  - $\eta - \eta'$ mixing implementation
  - Hybrid method (symmetry + quark dynamics), overcomplete basis: more flavour topologies than true $SU(3)_F$ invariants

Alternative flavour symmetry approaches

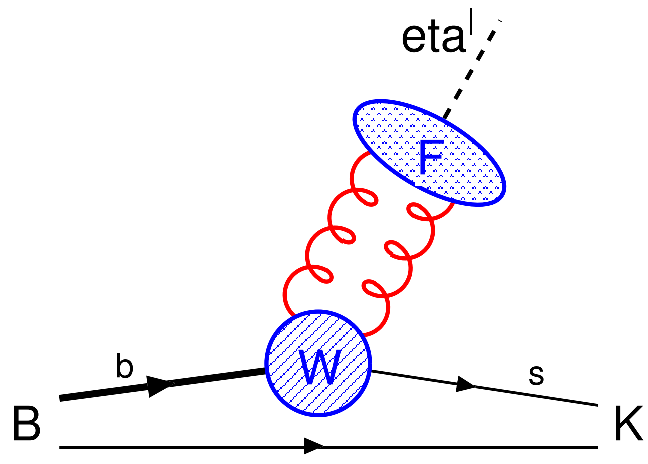
- Comparison of different $\eta - \eta'$ mixing implementations (single angle, two angles [Feldman, Kroll, Stech])
  $\longrightarrow$ results practicaly unchanged
- "Pure" $SU(3)_F$ symmetry approach (generalization of [Savage and Wise (1989)])

$$H_{\text{eff}} = a\, B_k H(3)^k P^j_i P^i_j + b\, B_i H(3)^k P^j_k P^i_j + c\, B_i H(\bar{6})^{ij}_k P^m_j P^k_m$$
$$+ d\, B_i H(15)^{ij}_k P^m_j P^k_m + e\, B_i H(15)^{jk}_m P^m_k P^i_j + \tilde{f}\, B_i H(3)^k P^i_k \eta_1$$
$$+ \tilde{g}\, B_i H(\bar{6})^{ij}_k P^k_j \eta_1 + \tilde{h}\, B_i H(15)^{ij}_k P^k_j \eta_1 + \tilde{s}\, B_k H(3)^k \eta_1 \eta_1$$

- we get $SP/P = 0.31 - 0.36$
- [Fu, He, Hsiao (2003)] $SP/P \approx 0.9$

Perturbative (dynamical) analysis

- [Atwood and Soni (1997)]
- [Halperin and Zhitnitsky (1997)]
- [Kagan and Petrov (1997)]
- [Hou and Tseng (1998)]
- [Datta, He and Pakvasa (1998)]
- [Du, Kim and Yang (1998)]
- [Ahmady, Kou and Sugamoto (1998)]
- [Ali, Chay, Greub and Ko (1998)]
- [Kou and Sanda (2002)]
- [Xiao, Chao and Li (2002)]
- [Beneke and Neubert (2002)]
- [Fritzsch and Zhou (2003)]



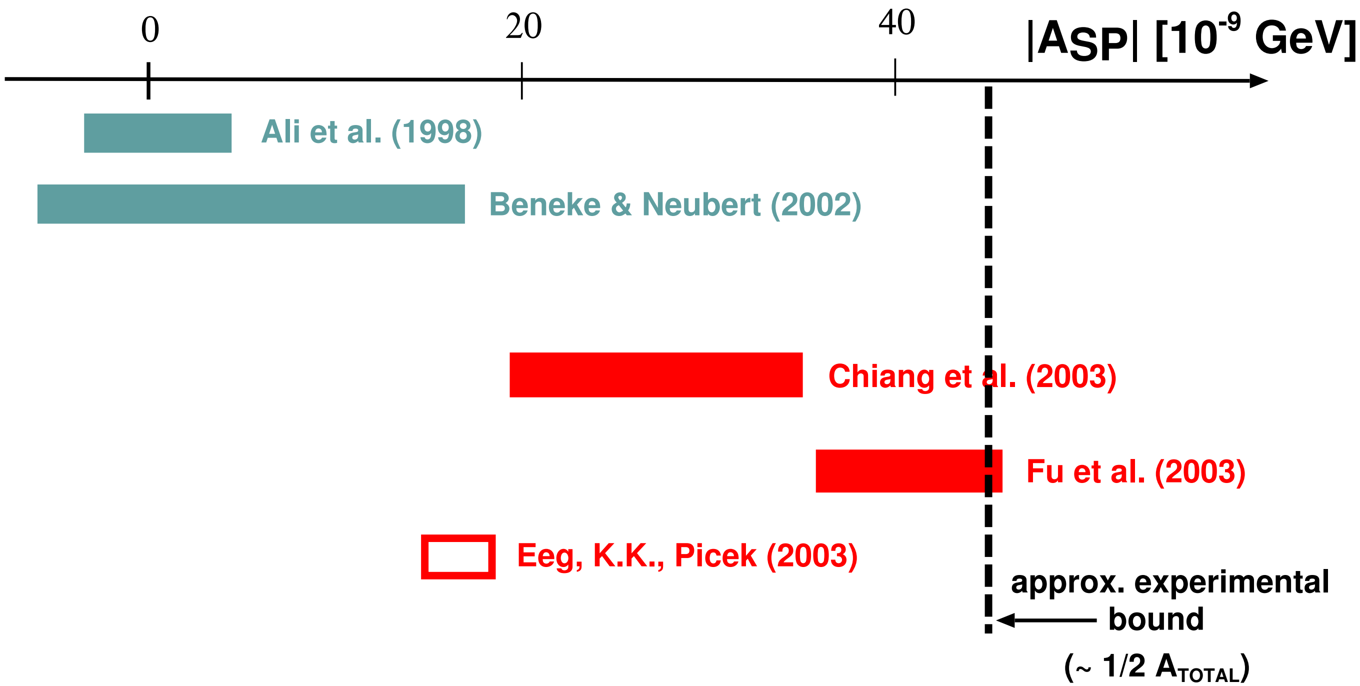
Generic features

- $b \to sgg$ transition described by either
  - $H_{\rm eff}^{\rm ew} = \frac{G_F}{\sqrt{2}} \sum C_i O_i \qquad O_1 = (\bar{u}b)_{\text{V-A}}(\bar{s}u)_{\text{V-A}}, \cdots$
  - or $\quad H_{\rm eff}(b \to sgg)$ [Simma and Wyler (1990)]
- $gg\eta'$ vertex described by

$$\langle \eta' | \frac{\alpha_s}{4\pi} G_{\mu\nu} \tilde{G}^{\mu\nu} | 0 \rangle = f_{\eta'} m_{\eta'}^2$$

- But:
  - This is appropriate for on-shell/soft gluons
  - $SP \propto \left(C_2 + \frac{C_1}{N_C} \right) = a_2 \simeq 0.2 \;\Rightarrow\; SP \ll P,\ \mathcal{A}_{\text{exp.}}$

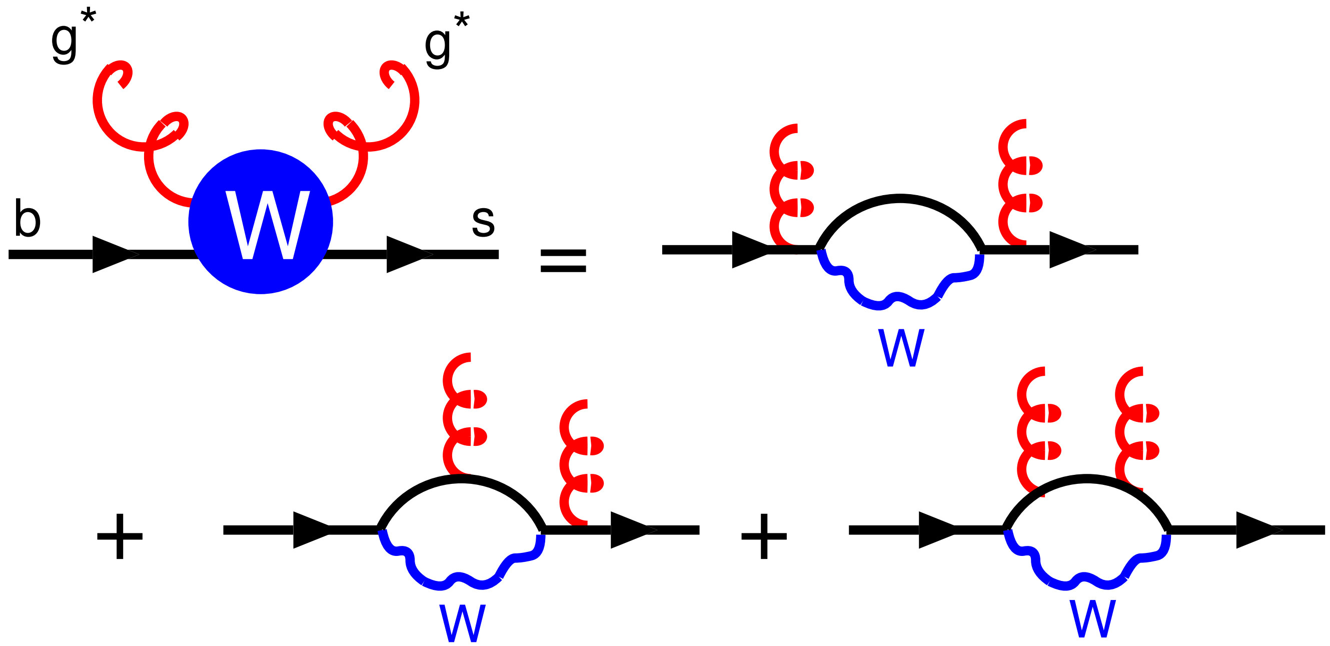
Comparison of two approaches I



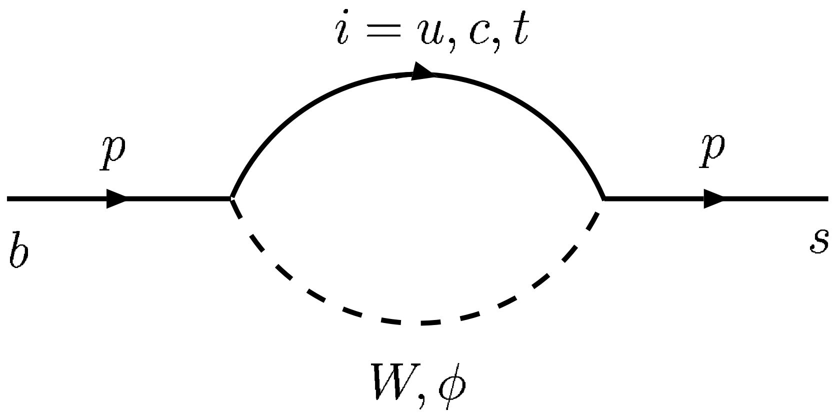
What about hard off-shell gluon contribution? Can it explain the discrepancy?

$b \longrightarrow sg^*g^*$ amplitude

- [Simma and Wyler (1990)]: small external momenta — $p_b, p_s, p_g \ll m_W$
- This work: $p_b, p_s \longrightarrow 0$, but general p_g
- Building blocks:



$b \longrightarrow s g^* g^*$ (self-energy)



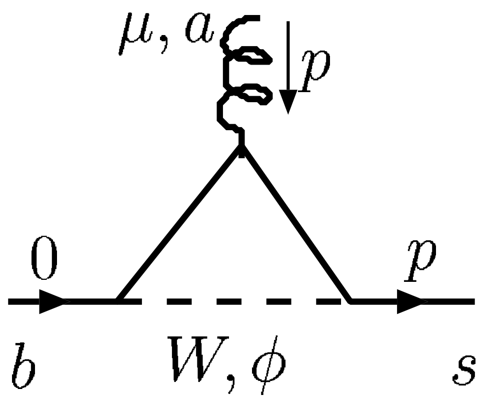
$$= \frac{i}{4\pi^2} \frac{G_F}{\sqrt{2}} \Sigma(p)$$

$$\boldsymbol{\Sigma(p)} = -M_W^2 \not{p} L - 2M_W^2 \left(1 + \frac{m_i^2}{2M_W^2}\right) \not{p} L \int_0^1 \mathrm{d}x (1-x) \ln \frac{D}{\mu_*^2}$$

$$- \int_0^1 \mathrm{d}x \left[(1-x) m_b m_s \not{p} R - m_i^2 (m_b R + m_s L)\right] \ln \frac{D}{\mu_*^2}$$

$$\ln \mu_*^2 = \tfrac{1}{\epsilon} - \gamma_E + \ln 4\pi\mu^2$$

$b \to s g^* g^*$ (Triangle)



$$= \frac{i}{4\pi^2}\frac{G_F}{\sqrt{2}} g_s t^a\, \Gamma^\mu(0,p,-p)$$

$$\boldsymbol{\Gamma^\mu(0,p,-p)} = \frac{4M_W^2}{m_i^2 - M_W^2}\left(1+\frac{m_i^2}{2M_W^2}\right)(p^2 g^{\mu\nu} - p^\mu p^\nu)\gamma_\nu L \int_0^1 \mathrm{d}x x(1-x)\ln\frac{D}{C}$$

$$+ M_W^2 \gamma^\mu L + 2M_W^2\left(1+\frac{m_i^2}{2M_W^2}\right)\gamma^\mu L \int_0^1 \mathrm{d}x(1-x)\ln\frac{D}{\mu_*^2}$$

$$D = x m_i^2 + (1-x)M_W^2 - x(1-x)p^2$$

$$C = m_i^2 - x(1-x)p^2$$

$b \longrightarrow s g^* g^*$ (Box)

$$I^{\mu\nu}(0,0,-p,p) = \frac{2M_W^2}{m_i^2 - M_W^2}\left(1 - \frac{m_i^2}{2M_W^2}\right)(-i\epsilon^{\mu\nu\rho\sigma} p_\sigma \gamma_\rho L) \times$$

$$\times \int_0^1 \mathrm{d}x (1-x) \left\{ (3x-1)\mathbb{Y}_1 + \left[x^2(1-x)p^2 + (x+1)m_i^2\right]\mathbb{Y}_2 \right\}$$

$$+ \frac{2M_W^2}{m_i^2 - M_W^2}\left(1 + \frac{m_i^2}{2M_W^2}\right) \int_0^1 \mathrm{d}x (1-x) \left\{ [-(x+1)\not{p} g^{\mu\nu} - (x-1)(p^\mu \gamma^\nu + p^\nu \gamma^\mu)]\mathbb{Y}_1 \right.$$

$$+ \left(x^2(1-x)\left[-(p^\mu\gamma^\nu + p^\nu\gamma^\mu)p^2 + \not{p}(4p^\mu p^\nu - g^{\mu\nu}p^2)\right] \right.$$

$$\left. \left. + [-(x+1)\not{p} g^{\mu\nu} - (x-1)(p^\mu\gamma^\nu + p^\nu\gamma^\mu)]m_i^2 \right) \mathbb{Y}_2 \right\} L$$

$\mathbb{Y}_{1,2}$ = complicated functions of x, m_i^2, M_W^2, p^2

$b \longrightarrow sg^*g^*$ (Complete)

$$\mathcal{A} = \mathrm{i}\frac{\alpha_s}{\pi}\frac{G_F}{\sqrt{2}}\bar{s}(0)t^b t^a \sum_i \lambda_i T_{i\mu\nu} b(0)\epsilon_a^{\mu}(-p)\epsilon_b^{\nu}(p) \; + \; \text{(crossed)} \; ,$$

$$T_i^{\mu\nu} = T_{i\text{Box}}^{\mu\nu} + T_{i\text{Triangle}}^{\mu\nu} + T_{i\text{Self-energy}}^{\mu\nu} \; .$$

$$T_i^{\mu\nu} = (-i\epsilon^{\mu\nu\rho\sigma} p_\sigma \gamma_\rho L)\, A_i \; + \; (\mu\text{-}\nu \text{ symmetric part})$$

$$A_i = -\frac{8M_W^2}{m_i^2 - M_W^2}\left(1 + \frac{m_i^2}{2M_W^2}\right)\int_0^1 \mathrm{d}x\, x(1-x)\ln\frac{D}{C}$$

$$+\frac{2M_W^2}{m_i^2 - M_W^2}\left(1 - \frac{m_i^2}{2M_W^2}\right)\int_0^1 \mathrm{d}x(1-x)\left\{(3x-1)\mathbb{Y}_1 + \left[x^2(1-x)p^2 + (x+1)m_i^2\right]\mathbb{Y}_2\right\}$$

$\eta' g^* g^*$ form-factor I

- $g^* g^* \eta'$ form-factor $F_{\eta' g^* g^*}$ poorly known
  $\longrightarrow$ recent improvements via perturbative QCD:
  - [Muta and Yang (2000)]
  - [Ali and Parkhomenko (2002,2003)]
  - [Kroll and Passek-Kumericki (2003)]
- $F_{\eta' g^* g^*}$ defined via $\eta' \to g^*(k_1) g^*(k_2)$ amplitude:

$$N^{ab}_{\mu\nu}(\bar{Q}^2, \omega) = -i F_{\eta' g^* g^*}(\bar{Q}^2, \omega)\, \epsilon_{\mu\nu k_1 k_2}\, \delta^{ab} \ ,$$

$$\bar{Q}^2 = -\frac{k_1^2 + k_2^2}{2} \qquad \omega = \frac{k_1^2 - k_2^2}{k_1^2 + k_2^2}$$

$\eta' g^* g^*$ form-factor II

- For $\bar{Q}^2 \gtrsim m_b^2$

$$F_{\eta' g^* g^*}(\bar{Q}^2, 0) = 4\pi\alpha_s(\bar{Q}^2)\frac{f^1_{\eta'}}{\sqrt{3}\bar{Q}^2}\left(1 - \underbrace{\frac{1}{12}B^g_2(\bar{Q}^2)}_{|\eta'\rangle=|gg\rangle}\right)$$

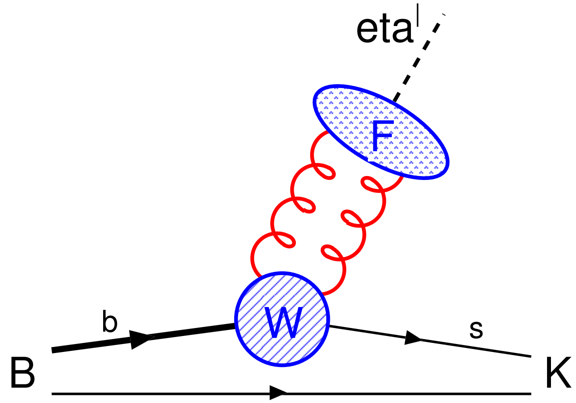
$$f^1_{\eta'} \approx 1.15\sqrt{2} f_\pi$$

- Double suppression of $F_{\eta' g^* g^*}$:

$$\left.\begin{array}{c} 1/\bar{Q}^2 \\ \alpha_s(\bar{Q}^2) \text{ running} \end{array}\right\} \quad \text{for } \bar{Q}^2 \gg$$

Gluing two pieces together

- Combining amplitudes for $\boldsymbol{b \to s g^* g^*}$ and $\boldsymbol{g^* g^* \to \eta'}$

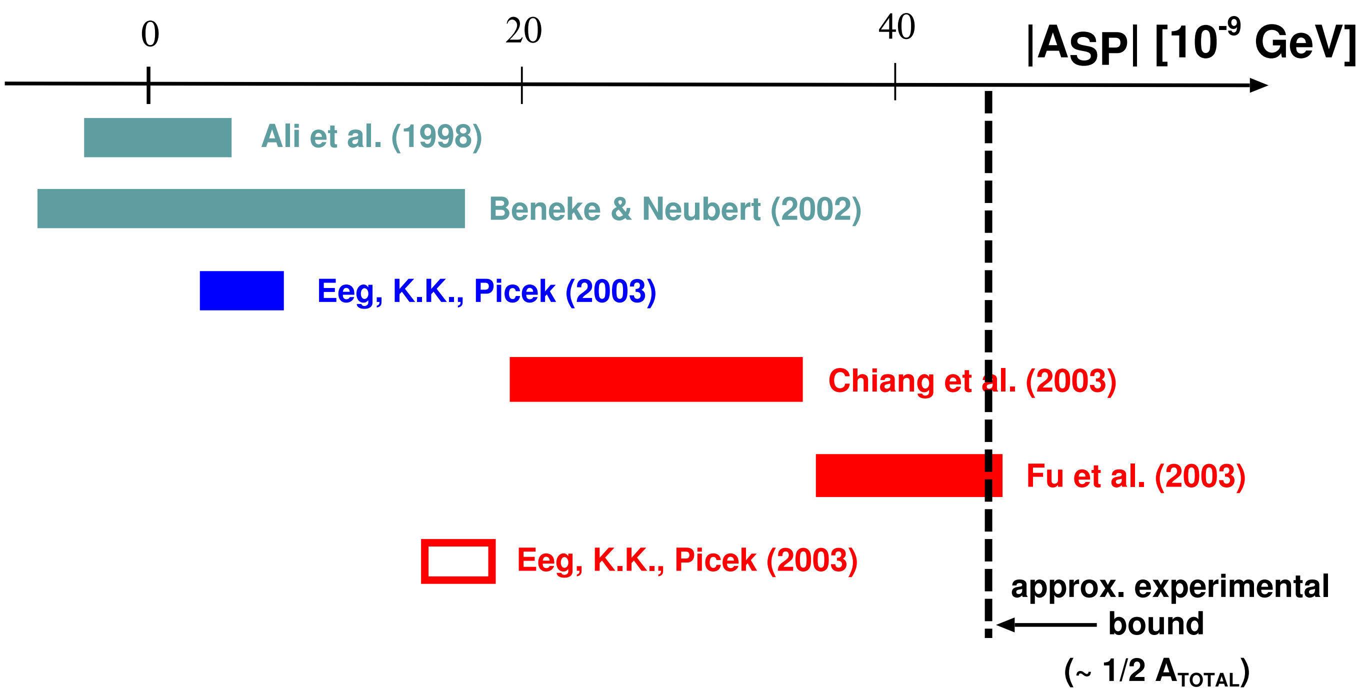


$$\mathcal{A}(b \to s\eta') = \frac{G_F}{8\sqrt{2}\pi^3} \left(\phi_{\eta'} \bar{s} \not{P}_{\eta'} L b\right) \sum_{i=u,c,t} \lambda_i$$

$$\times \int_{\mu^2 \sim m_b^2}^{M_W^2} dQ^2\, \alpha_s(Q^2)\, F_{\eta' g^* g^*}(Q^2)\, A_i(-Q^2)$$

- $\boldsymbol{\mathcal{A}(b \to s\eta') \;\to\; \mathcal{A}(B \to K\eta')}$ via factorization

Comparison of two approaches II



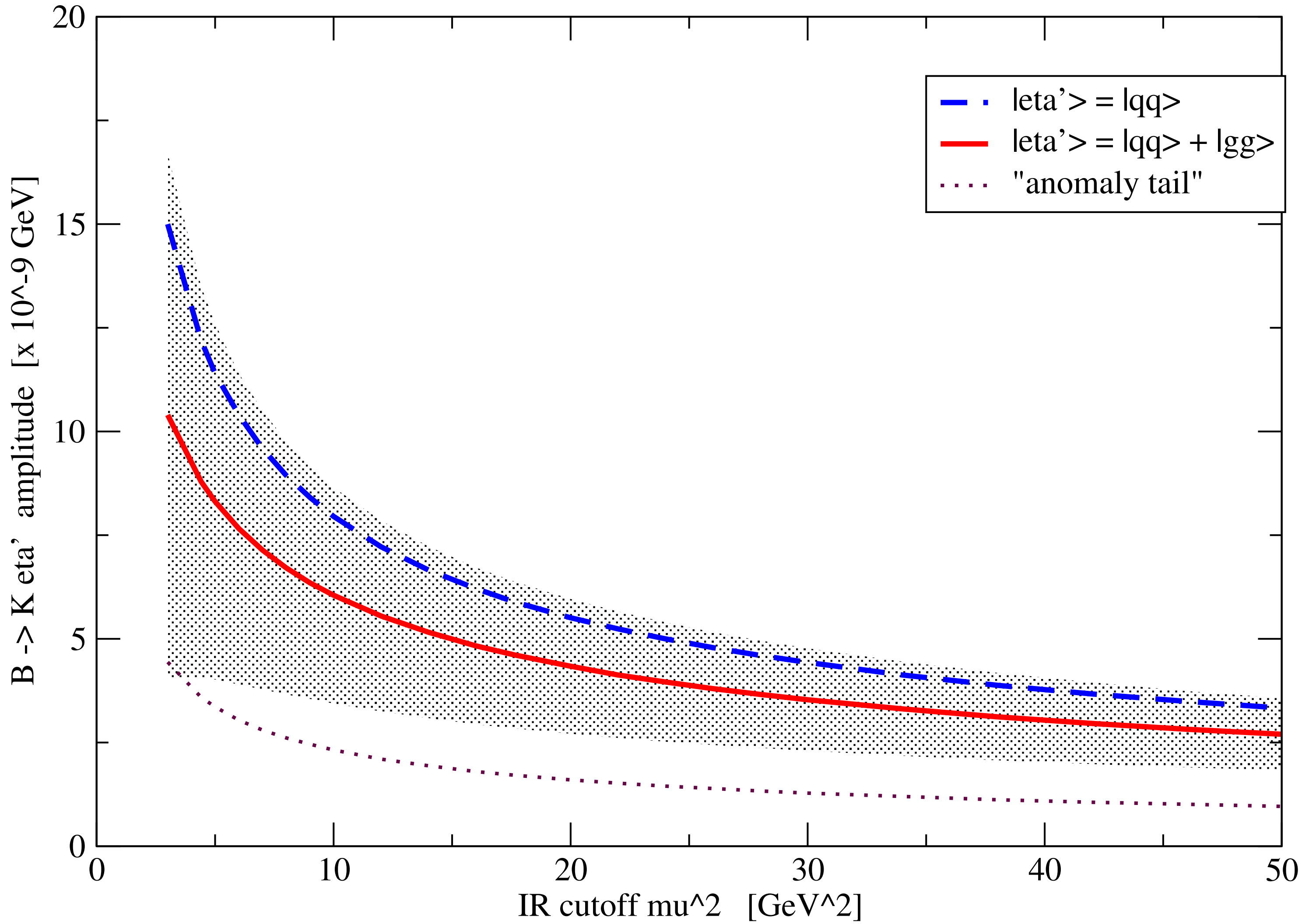
- (One must add SD (blue) on top of LD (gray-blue) and than compare with SU(3) (red).)
- Discrepancy smaller but still exists!

Conclusions

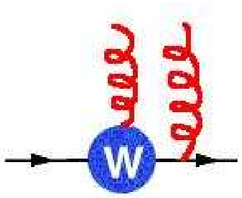
- Singlet penguin gluonic mechanism has significant but not dominant role in $B \to K\eta'$ amplitude
- Maybe "significant" $\longrightarrow$ "detectable"
- No new physics is needed, but better understanding of the discrepancy between two described approaches would be welcome

The End

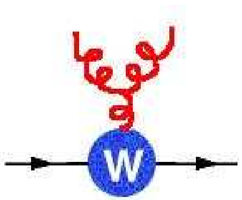
IR cut-off dependence

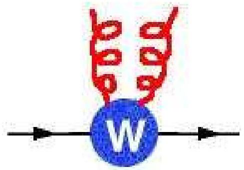


F1-F2 interplay

- 

$\sim F_1(x)(p^2\gamma^\mu - \not{p}p^\mu)L - F_2(x)i\sigma_{\mu\nu}q^\nu m_b R$

- 

$\sim$ (F_1-terms) + (F_2-terms)


$\sim$ (F_1-terms) + (F_2-terms)

- $x \ll 1 \;\Rightarrow\; (F_1 \sim \ln x) \;\gg\; (F_2 \sim x^2 \ln x)$

- F_1 terms cancel for on-shell or soft gluons (Ward identities, low-energy theorem [Low(1958)]) $\Rightarrow$ suppression

- but not for hard off-shell gluons ([Witten (1977)])!

Hiperindex