

New opportunities for kaonic atoms measurements from CdZnTe detectors

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072 Abstract

073 We present the tests performed by the SIDDHARTA-2 collaboration at
074 the **DAΦNE** collider with a quasi-hemispherical CdZnTe detector. The
075 very good room-temperature energy resolution and efficiency in a wide
076 energy range show that this detector technology is ideal for studying
077 radiative transitions in intermediate and heavy mass kaonic atoms. The
078 CdZnTe detector was installed for the first time in an accelerator envi-
079 ronment to perform tests on the background rejection capabilities, which
080 were achieved by exploiting the SIDDHARTA-2 Luminosity Monitor. A
081 spectrum with an ²⁴¹Am source has been acquired, with beams circulat-
082 ing in the main rings, and peak resolutions of 6% at 60 keV and of 2.2% at
083 511 keV have been achieved. The background suppression factor, which
084 turned out to be of the order of $\simeq 10^{-6} - 10^{-7}$, opens the possibility
085 to plan for future kaonic atom measurements with CdZnTe detectors.

086 **Keywords:** CdZnTe detectors, Kaonic Atoms, Offline Trigger

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1 Introduction

Kaonic atoms are formed when a K^- is moderated inside a target until it reaches a low enough kinetic energy to be stopped, replacing one of the outer electrons and forming an exotic atom in a highly excited state. The kaonic atom then undergoes atomic cascade to the ground state. These systems provide an ideal tool to study the low-energy regime of Quantum Chromodynamics (QCD) since, due to the much heavier K^- mass with respect to the e^- one, the lower levels are close enough to the nucleus to be influenced by the short-range strong interaction between the nucleus and the K^- [1].

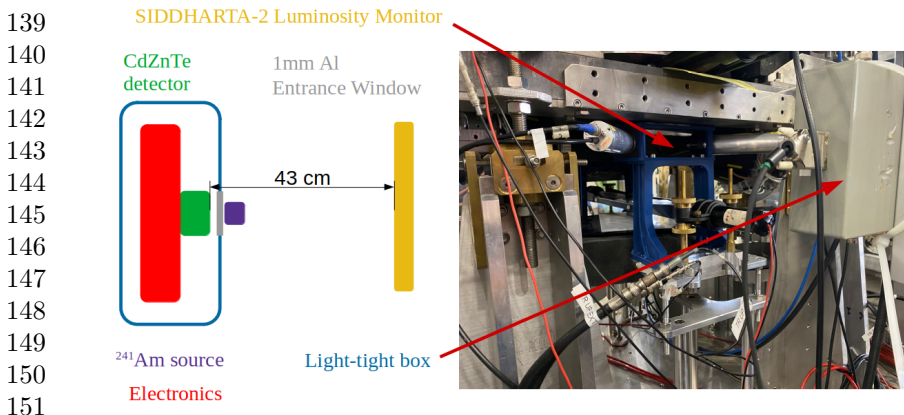
Kaonic atoms have been intensively studied in the 1970s and 1980s with a series of measurements, still representing today the main database for low-energy antikaon-nucleon studies [2–7]. More recently, a second generation of experiments performed new and more accurate measurements on kaonic hydrogen and helium [8–13], whose results were in better agreement with the theoretical predictions but not compatible with the old experiments [14].

The technological boost in radiation detectors in the last decade stimulated the strangeness nuclear physics community to plan for new experiments aiming to renew the database on kaonic atoms, in particular at the DAΦNE collider at the INFN National Laboratories of Frascati [15]. To detect the radiation emitted from the various transitions, a special role could be played by cadmium–zinc–telluride (CdZnTe) detectors, thanks to their excellent performances at room temperatures in terms of energy resolution (few % FWHM), linearity and fast readout (few tens of ns) in a broad energy region ranging from few keV to some MeV [16, 17].

One of the most challenging difficulties to be overcome in a facility like DAΦNE is the high electromagnetic and hadronic background to which these detectors will be exposed, since their small active area requires to place them as close as possible to the Interaction Point (IP), to maximize the solid angle acceptance. In this work we present the first tests ever performed with such types of detectors in an accelerator; they were specifically aimed to assess the performances in the presence of a high machine background, similarly to what was already performed for the Silicon Drift Detectors by the SIDDHARTA-2 collaboration [18]. Their fast timing response allows for a high-effective background suppression, achieved by using a trigger system based on kaons detection via Time of Flight (ToF), crucial for the possibility to observe the radiation emitted from the kaonic atom transitions [19].

2 Experimental setup

In June 2022, a first prototype of a quasi-hemispherical CdZnTe detector system, with an active surface of 1 cm^2 and a thickness of 5 mm, was installed near the IP of the DAΦNE collider, vertically aligned with the SIDDHARTA-2 Luminosity Monitor (LM) [19], at a distance of 43 cm. A (not in scale) top view of the geometry of the setup and a picture of the prototype as installed at DAΦNE are shown in Fig. 1.



152 **Fig. 1** *Left:* Top view of the geometry of the setup. *Right:* Picture of the light-tight box
153 containing the CdZnTe prototype installed at the DAΦNE collider.

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155 The prototype, manufactured by the IMEM-CNR of Parma, was enclosed in
156 a light-tight box with a 1 mm thick aluminum entrance window, matching
157 the detector's active surface, on top of which an ²⁴¹Am radioactive source was
158 placed, producing a $\simeq 500$ Hz signal in the CdZnTe detector.

159 The detector was connected to low-noise (equivalent noise charge ENC of 100
160 electrons) charge-sensitive preamplifiers (CSPs) and processed by 8-channel
161 digital electronics. Both the CSPs and the digital electronics were developed
162 at DiFC of University of Palermo (Italy). The digital electronics consisted of
163 a CAEN DT5724 digitizer driven by an original firmware [20–22].

164 The signals from the LM were also acquired, with the same digitizer, to perform
165 an offline data selection (trigger) when a charged kaon pair was produced in the
166 horizontal plane. The digital signals from the LM were processed by an ORTEC
167 566 Time-to-Amplitude Converter (TAC) module. The LM is based on the J-
168 PET technology developed for the imaging of electron-positron annihilation
169 [23, 24]. A detailed description of the LM and of the discrimination between
170 the kaons and the Minimum Ionizing Particles (MIPs) produced during the
171 e^+e^- collisions can be found in [19].

172 The data reported in this paper have been acquired for a total of 72 hours,
173 during which the accelerator delivered e^- and e^+ beams with average currents
174 of $\simeq 500$ mA and $\simeq 270$ mA, respectively.

176 3 Results

177
178 The spectrum acquired in the 72 hours test with the CdZnTe is plotted,
179 without any data selection requirement, in Fig. 2.

180 The spectrum was obtained without any selection. The ²⁴¹Am peaks are visible
181 together with the Neptunium (Np) ones, both coming from the source decay
182 chain. The Lead (Pb) X-rays come from the activation of the SIDDHARTA-2
183 setup lead shielding by the MIPs lost from the beams. The three escape peaks
184 of the main ²⁴¹Am peak, due to Cd, Zn and Te, are also observed [20]. The

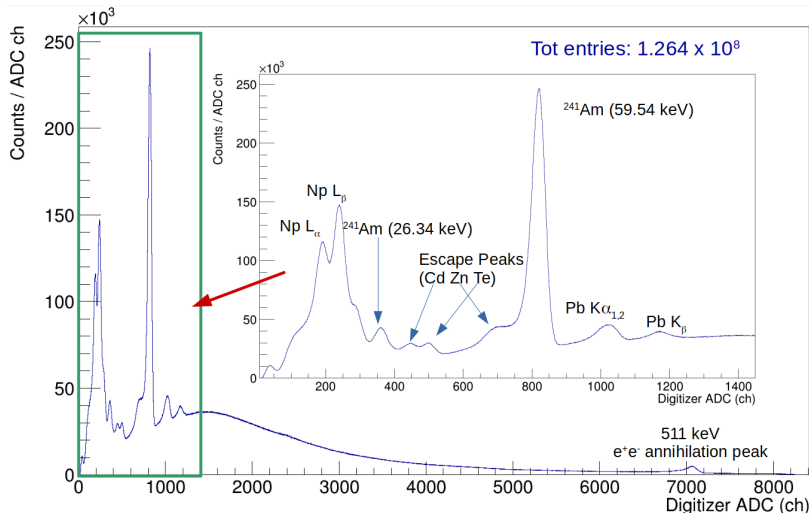


Fig. 2 CdZnTe spectrum acquired in 72 hours at DAΦNE without any trigger requirement.

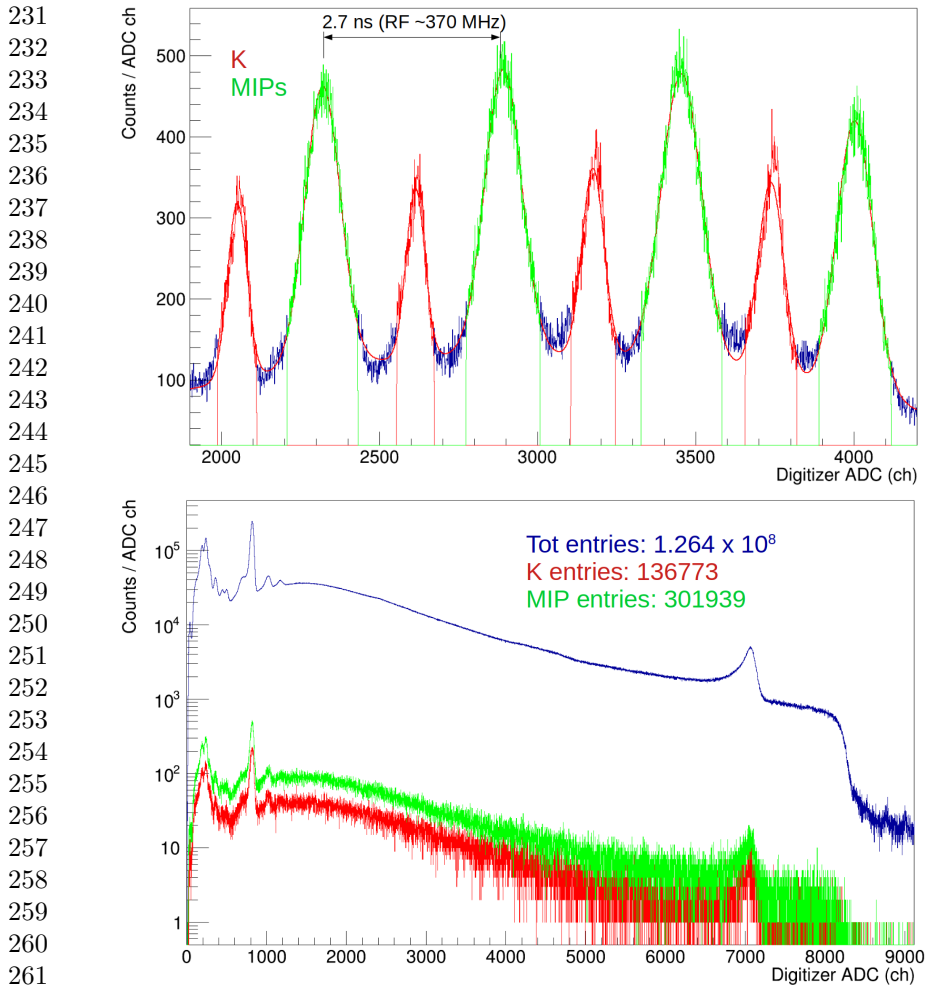
measured peak resolutions are of 6% at 60 keV and of 2.2% at 511 keV. To simulate the rejection factor of the background of the DAΦNE machine, the TAC signals have been used to select only events in which a K^+K^- pair passed through the two scintillators of the LM. In this way, the case in which only events correlated in time with the formation of kaonic atoms in a target placed in front of the CdZnTe entrance window are retained, is simulated.

In the upper part of Fig. 3, the TAC spectrum from the SIDDHARTA-2 LM is shown, where red and green peaks correspond to kaons and MIPs, respectively. The two LM scintillators are placed at a specific distance, with respect to the IP, such as to have a difference in the arrival time between kaons and MIPs producing two clear peaks, visible in the figure. The ~ 370 MHz radiofrequency (RF) of the DAΦNE collider is used in coincidence with the signals from the LM to get a few ps resolution timestamp for each event. To provide a reference for each collision, a Constant Fraction Discriminator (CFD) must be employed. The CFDs are usually limited to work at 200 MHz; to overcome this limitation, the RF/4 is then used, at a frequency of ~ 90 MHz. As a consequence, every coincidence event in the LM discriminators can be randomly associated in time with one of the four collisions; this is reflected as the four double structures of Fig. 3.

The events from the radioactive source are uncorrelated in time with the DAΦNE collisions; to test the background suppression, we thus require that a signal from the LM occurs just before a signal in the CdZnTe detector, simulating the case where X-rays from the transitions of kaonic atoms hit the detector immediately after the detection of a kaon pair in the LM.

The lower part of Fig. 3 shows the spectrum with no time selection from the TAC (blue) overimposed with the spectra obtained requiring the presence either of a kaon (red) or a MIP (green) previous to the CdZnTe signal. The

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262 **Fig. 3** *Top:* TAC spectrum from the SIDDHARTA-2 LM; red and green peaks correspond
 263 to kaons and MIPs, respectively. *Bottom:* CdZnTe spectrum with no time selection from the
 264 TAC (blue) overimposed on the spectra obtained in coincidence with kaons (red) or MIPs
 265 (green) events.

266 shapes of these spectra are all identical and reflect the absence, in this test, of
 267 a target where kaonic atoms can be formed. The trigger-induced suppression
 268 consists of the rejection of all those events in which a detected decay from the
 269 ^{241}Am source is not (randomly) occurring in coincidence with the signal on
 270 the LM. The requirement of a kaon or a MIP signal within two CdZnTe events
 271 induces a suppression of the background of a factor $\simeq 10^3$.

272 The events from the source occur randomly in time with respect to the col-
 273 lisions and, consequently, to the LM signals. The distributions of the time
 274 difference, ΔT , between a photon detection by the CdZnTe detector and a
 275 signal in the SIDDHARTA-2 LM are shown in Fig. 4 for kaons (red), MIPs
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(green) and all the other events out of the K/MIP selections (blue). For each event on the LM, processed with the TAC and acquired with the digitizer, the distributions show the ΔT s with respect to the subsequent event on the CdZnTe detector.

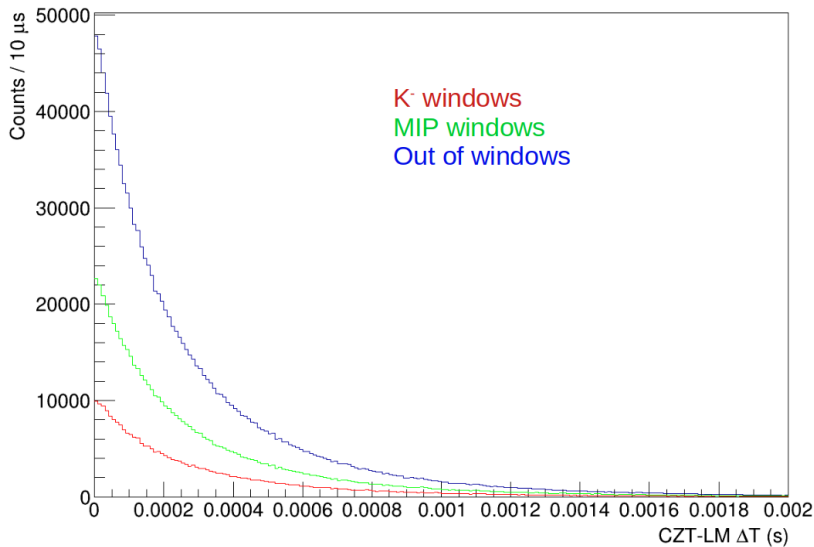


Fig. 4 Distribution of the time difference between a photon detection by the CdZnTe detector and a signal in the SIDDHARTA-2 LM.

From Fig. 4 it turns out that ^{241}Am events occur with a time difference up to 2 ms, which is exactly what expected from the 500 Hz rate of the ^{241}Am source. When operated with a target to perform kaonic atoms measurements, the ΔT distribution is expected to have two components: a flat background, collecting all the non-kaonic events on the CdZnTe occurring randomly in time with respect to the LM ones, and a second one where, on the contrary, the events correlated with a K^- passing through the LM, stopping in the target and forming the kaonic atoms will produce a peak.

In a similar way to what was observed for the Silicon Drift Detectors (SDDs) of the SIDDHARTA-2 setup [18], this peak will have a FWHM reflecting both the timescale of the physical processes described above and that of the CdZnTe readout. In [18], this FWHM is ranging from 430 to 950 ns, depending on the temperature of the SDDs, but the faster readout of CdZnTe detectors suggests that smaller values could be achieved. In Tab. 1, we report the number of events on the CdZnTe satisfying the request to have a coincidence with a K^- in the LM (K_{TAC}^-) without any ΔT request as well as with additional requirements of $\Delta T < 1000, 500, 300, 100$ ns. The resulting background suppression factors are also reported.

Depending on the time window, suppression factors of the order of $\simeq 10^{-6} - 10^{-7}$ are found. In particular, the request to have a TAC signal in

Request	Events	Rejection factor
No request	1.26×10^8	1
K_{TAC}^-	136359	1.1×10^{-3}
$K_{\text{TAC}}^-, \Delta T < 1 \mu\text{s}$	1096	8.7×10^{-6}
$K_{\text{TAC}}^-, \Delta T < 500 \text{ ns}$	605	4.8×10^{-6}
$K_{\text{TAC}}^-, \Delta T < 300 \text{ ns}$	374	3.0×10^{-6}
$K_{\text{TAC}}^-, \Delta T < 100 \text{ ns}$	124	9.8×10^{-7}

Table 1 Number of events on the CdZnTe satisfying the request to have a coincidence with a K^- in the LM (TAC), without any ΔT request as well as with additional requirements of $\Delta T < 1000, 500, 300, 100 \text{ ns}$, and the resulting background suppression factors.

time between two CdZnTe ones is responsible for the major suppression of the asynchronous background (mainly due to MIPs lost from the beams), while further reductions of the coincidence time window not only contribute to this suppression, but also, requiring a K^+K^- pair in the LM, help in the rejection of the synchronous component of the background [18].

4 Conclusions

In this work we presented the first tests performed with a quasi-hemispherical CdZnTe detector in an accelerator environment like the DAΦNE collider at the INFN Laboratories of Frascati, where an ^{241}Am spectrum was acquired for 72 hours, with the beams circulating in the main rings, showing peak resolutions of 6% at 60 keV and of 2.2% at 511 keV. We successfully measured ad suppression of the machine background with a trigger system, exploiting the fast readout capability of the device. The $\simeq 10^{-6} - 10^{-7}$ obtained rejection factors represent a very promising result which encourage the SIDDHARTA-2 collaboration to include, in its future data-taking, measurements of radiative transitions from several intermediate and high mass kaonic atoms to be performed with CdZnTe detectors.

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